

CO₂ bubble generation and migration during magma–carbonate interaction

L. S. Blythe^{1,2} · F. M. Deegan^{1,3} · C. Freda⁴ · E. M. Jolis¹ · M. Masotta⁵ · V. Misiti⁴ · J. Taddeucci⁴ · V. R. Troll^{1,4}

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Abstract We conducted quantitative textural analysis of vesicles in high temperature and pressure carbonate assimilation experiments (1200 °C, 0.5 GPa) to investigate CO₂ generation and subsequent bubble migration from carbonate into magma. We employed Mt. Merapi (Indonesia) and Mt. Vesuvius (Italy) compositions as magmatic starting materials and present three experimental series using (1) a dry basaltic-andesite, (2) a hydrous basaltic-andesite (2 wt% H₂O), and (3) a hydrous shoshonite (2 wt% H₂O). The duration of the experiments was varied from 0 to 300 s, and carbonate assimilation produced a CO₂-rich fluid and CaO-enriched melts in all cases. The rate of carbonate assimilation, however, changed as a function of melt viscosity, which affected the 2D vesicle number, vesicle volume, and vesicle size distribution within each experiment. Relatively low-viscosity melts (i.e. Vesuvius experiments) facilitated efficient removal of bubbles from the reaction site. This allowed carbonate assimilation to continue unhindered and large volumes of CO₂ to be

liberated, a scenario thought to fuel sustained CO₂-driven eruptions at the surface. Conversely, at higher viscosity (i.e. Merapi experiments), bubble migration became progressively inhibited and bubble concentration at the reaction site caused localised volatile over-pressure that can eventually trigger short-lived explosive outbursts. Melt viscosity therefore exerts a fundamental control on carbonate assimilation rates and, by consequence, the style of CO₂-fuelled eruptions.

Keywords CO₂ · Carbonate assimilation · Melt viscosity · Bubble size distribution · Eruption style

Introduction

Earth's long-term volatile cycle reflects mass exchange between the deep Earth and the atmosphere, with volcanoes representing the major agents of volatile transfer. When magma migrates through crustal rocks *en route* from the mantle to the surface, it interacts thermally, chemically, and mechanically with the surrounding rocks. If these rocks are volatile rich, such as limestone and dolostone which contain structurally bound CO₂, magma–crust interaction can liberate crustal volatiles into the magmatic system, where they will eventually be released to the atmosphere via eruption or degassing.

While long-term magmatic degassing constantly delivers volatiles to the atmosphere, magma–crust interaction, in turn, may drive fluctuating peaks of “extra” CO₂ sourced from the crust. The latter process has been recognised in the recent volcanic gas record at Merapi volcano, Indonesia (e.g. Troll et al. 2012) and at Vesuvius volcano, Italy (Iacono Marziano et al. 2009) and is thought to have played a major role in perturbations to the global carbon cycle in

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✉ V. R. Troll
valentin.troll@geo.uu.se

¹ Department of Earth Sciences, Centre for Experimental Mineralogy, Petrology, and Geochemistry (CEMPEG), Uppsala University, Villavägen 16, 752 36 Uppsala, Sweden

² School of Physical and Geographical Science, Keele University, Keele, UK

³ Department of Geological Sciences, Stockholm University, Stockholm, Sweden

⁴ Istituto Nazionale di Geofisica e Vulcanologia (INGV), Rome, Italy

⁵ Bayerisches Geoinstitut, Universität Bayreuth, Bayreuth, Germany

Earth's past (Ganino and Arndt 2009; Svensen et al. 2009; Johnston et al. 2011; Lee et al. 2013).

Indeed, it has been suggested that CO₂ emissions were several times greater than at present in, e.g. the Cretaceous, which, by consequence, had a warmer climate (Johnston et al. 2011; Lee et al. 2013). Excess CO₂ outgassing in the Cretaceous notably cannot be explained by pure decarbonation of subducted materials alone, especially as this process is thought to be rather inefficient based on thermodynamic models (Kerrick and Connolly 2001). Although less pronounced than in some earlier periods of Earth's history, an imbalance still exists between the amount of CO₂ released during subduction and the amount of CO₂ emitted at arc volcanoes at present (Dasgupta 2013). Carbonate-rich upper crustal lithologies beneath arcs have therefore been considered a missing piece in the carbon cycle puzzle (cf. Johnston et al. 2011; Troll et al. 2012; Dasgupta 2013; Lee et al. 2013). In addition to the repercussions for the global carbon cycle on a geological timescale, magma–carbonate interaction can also contribute to driving short-lived explosive eruptions at hazardous volcanoes besides Vesuvius and Merapi, including sites such as Popocatepetl, Mexico, and Yellowstone, USA (Goff et al. 2001; Werner and Brantley 2003; Schaaf et al. 2005; Chadwick et al. 2007; Iacono Marziano et al. 2009; Deegan et al. 2010, 2011; Dallai et al. 2011; Troll et al. 2012, 2013; Borisova et al. 2013). In this context, carbonate assimilation causes the host magma to become progressively more potassic and silica undersaturated, while the simultaneously produced CO₂-rich fluid phase is less well understood (e.g. Watkinson and Wyllie 1964, 1969; Freda et al. 2008; Gaeta et al. 2009; Mollo et al. 2010). Therefore, the dynamic aspects of CO₂ bubble nucleation, growth, and transport during carbonate assimilation and their combined effects on volcanic activity have still to be fully resolved.

To investigate the textural and chemical evolution of CO₂ bubbles during carbonate assimilation, time-constrained, high pressure–temperature experiments at conditions appropriate for mid- to upper-crustal magma reservoirs have been carried out. Starting materials from Merapi and Vesuvius volcanoes were doped with carbonate fragments and demonstrated both rapid production of CaO-enriched melts and CO₂ vesiculation in the vicinity of the added carbonate (Deegan et al. 2010, 2011; Jolis et al. 2013). In order to better comprehend the overall dynamics of the volcanic system during CO₂-fuelled eruptions and to further assess the role of magma–crust interaction in the subduction carbon cycle, it is imperative to quantify how CO₂-rich fluids are produced and mobilised in crustal systems (cf. Freda et al. 1997, 2011; Johnston et al. 2011; Troll et al. 2012; Dasgupta 2013; Jolis et al. 2013; Lee et al. 2013). We carried out analyses of vesicle

diameter, vesicle number, vesicle volume, and vesicle size distribution on previously characterised time-monitored carbonate assimilation experiments for Merapi and Vesuvius volcanoes (Deegan et al. 2010; Jolis et al. 2013) and compared these experimental sets directly for the first time. We then present a quantitative assessment of the controls on vesicle generation and mobility during crustal carbonate assimilation.

Geological background

Merapi and Vesuvius volcanoes are characterised by carbonate-hosted magmatic plumbing systems, wherein magma–carbonate interaction has been identified through isotopic studies of erupted minerals (e.g. feldspar, pyroxene, and olivine) and fumarolic CO₂ (e.g. Chadwick et al. 2007; Iacono Marziano et al. 2009; Dallai et al. 2011; Troll et al. 2012, 2013; Borisova et al. 2013). Notably, both volcanoes also have a history of explosive behaviour and are in proximity to major population centres, i.e. Yogyakarta in central Java (3.5 M inhabitants) and Naples in central Italy (>4 M inhabitants), making Merapi and Vesuvius significant threats to modern society.

Merapi is situated on the Eurasian plate, beneath which the Indian Ocean plate subducts at a rate of 6–7 cm/year (Tregoning et al. 1994). In central Java, crustal thickness is ~25 km (Curry et al. 1977) and Merapi sits at the margin of the Kendeng basin, which contains ≥10 km of upper crustal sedimentary strata (van Bemmelen 1949; Smyth et al. 2005). Tertiary marine strata, carbonates, volcanics, and volcanoclastics, probably similar to those outcropping in the nearby Djiwo hills, underlie Merapi volcano (van Bemmelen 1949; Hamilton 1979; Chadwick et al. 2007). The Merapi plumbing system consists of a plexus of magma pockets dispersed throughout the crust (Chadwick et al. 2013; Troll et al. 2013; van der Zwan et al. 2013), and calc-silicate xenoliths are frequent amongst the eruptive products, providing direct petrological evidence for magma–carbonate interaction in the mid- to upper-parts of the Merapi supply system (Camus et al. 2000; Gertisser and Keller 2003; Chadwick et al. 2007; Deegan et al. 2010; Troll et al. 2013, 2015). Fumarole chemistry indicates that carbonate assimilation is ongoing (Allard 1983; Troll et al. 2012), and what is more, post-earthquake fumarole emissions from the summit crater sampled in 2006 demonstrate exacerbated assimilation, likely as a consequence of crustal basement rupture (Deegan et al. 2011; Troll et al. 2012). Notably, the highly explosive 2010 Merapi eruption lacks abundant carbonate and skarn xenoliths, but shows strong geochemical evidence for intense carbonate digestion (Borisova et al. 2013). This “100-year” event was associated

with an enormous CO₂ output during the eruption (Surono et al. 2012), thus implying very efficient carbonate assimilation on that occasion.

Vesuvius is located in the Campania Province, Italy, where carbonate rocks extend from ≥ 2 km to at least 8 km depth (Zollo et al. 1996; Bruno et al. 1998; Auger et al. 2001; Peccerillo 2005). The current magma plumbing system is distributed over small chambers throughout much of the crust, e.g. between 8 and 22 km (Auger et al. 2001, Civetta et al. 2004). Geophysical data place the top of a magma chamber, which is of considerable horizontal extent, at ~ 8 km depth, i.e. at the base of carbonate crustal strata (Auger et al. 2001). A shallower (ca. 2–4 km) magma chamber fed older Plinian eruptions (e.g. Pompeii, AD 79; Avellino, 3500 BP), making shallow crustal magma–carbonate interaction a distinct possibility (Barberi et al. 1981; Civetta et al. 2004; Scaillet et al. 2008). Skarn xenoliths are a common feature in Vesuvius erupted products also and provide independent petrological evidence for magma–carbonate interaction (Fulignati et al. 2000, 2004a, b; Del Moro et al. 2001; Iacono Marziano et al. 2009; Dallai et al. 2011). The addition of volatiles sourced from carbonate must therefore play a significant role in magma chamber dynamics, especially since the current magma residence depth is assumed to be 3–4 km, e.g. for the 1944 eruption (cf. Fulignati et al. 2004b; Scaillet et al. 2008). In line with this inference, present-day fumarole emissions appear to record a considerable carbonate-derived CO₂ component (Graham et al. 1993; Martelli et al. 2004; Iacono Marziano et al. 2009).

Methodology

Experimental methods and analytical techniques

Piston cylinder experiments were previously carried out at the High Pressure High Temperature Laboratory of Experimental Volcanology and Geophysics at the Istituto Nazionale di Geofisica e Vulcanologia (INGV) in Rome, Italy (Deegan et al. 2010; Jolis et al. 2013). The experiments consisted of a carbonate fragment (ca. 6–10 mg) enclosed in pre-fused, powdered magmatic starting materials (ca. 24–42 mg) that were placed together in a platinum capsule. In all charges, the carbonate clast and magmatic material were ~ 20 and 80 % of the total volume in the capsule, respectively (Table 1). Analyses of all starting materials are reported in Table 2. The experiments were subsequently pressurised at room temperature to 0.5 GPa and then heated in two stages, from ambient temperature to 1180 °C at a rate of 200 °C/min, and at 20 °C/min until the experimental temperature was reached. Three experimental series are compared here: (1) a dry Merapi basaltic-andesite, (2) a hydrous Merapi basaltic-andesite (2 wt% H₂O), and (3) a hydrous Vesuvius shoshonite (2 wt% H₂O). Note that a total of 2 wt% H₂O was added to the hydrous experiments to account for the likely range of hydrous subduction zone compositions. The starting materials were then allowed to react at the target conditions for 0, 60, 90, and 300 s. The experiments were terminated by shutting down the power source, which caused quenching at a rate of ~ 2000 °C/min and allowed the textural interplay between limestone and

Table 1 Experimental conditions used for each series

Run ^{a,b}	Starting material	H ₂ O (wt%)	Carbonate (wt%) ^c	<i>T</i> (°C)	<i>P</i> (MPa)	<i>t</i> (s)
<i>Merapi dry series^a</i>						
379-16	Basaltic-andesite	0	19.6	1200	500	0
387-20	Basaltic-andesite	0	20.0	1200	500	60
376-10	Basaltic-andesite	0	19.0	1200	500	90
386-18	Basaltic-andesite	0	19.4	1200	500	300
<i>Merapi wet series^a</i>						
379-17	Basaltic-andesite	2.23	19.7	1200	500	0
387-21	Basaltic-andesite	2.23	19.3	1200	500	60
376-11	Basaltic-andesite	2.23	17.5	1200	500	90
386-19	Basaltic-andesite	2.23	19.3	1200	500	300
<i>Vesuvius wet series^b</i>						
443-V5	Shoshonite	2.02	17.1	1200	500	0
442-V3	Shoshonite	2.02	17.9	1200	500	60
441-V1	Shoshonite	2.02	21.9	1200	500	90
445-V7	Shoshonite	2.02	21.7	1200	500	300

^a Deegan et al. (2010)

^b Jolis et al. (2013)

^c Carbonate (wt%) = Assimilant weight (mg)/[Assimilant weight (mg) + Glass weight (mg)]

Table 2 Chemical compositions of starting materials

Oxide (wt%)	Average composition of basaltic-andesite glass (anhydrous) ^a	1 σ (10) ^a	Average composition of basaltic-andesite glass (hydrated) ^a	1 σ (9) ^{a,d}	Average composition of shoshonite glass (hydrated) ^b	1 σ (9) ^{b,d}	Merapi limestone ^a	Vesuvius limestone ^b
SiO ₂	54.11	0.60	51.83	0.43	49.85	0.21	0.28	0.02
TiO ₂	0.85	0.09	0.89	0.05	1.03	0.08	0.01	–
Al ₂ O ₃	18.98	0.17	18.08	0.24	15.95	0.67	0.13	–
FeO ^c	7.89	0.56	8.17	0.16	7.98	0.25	0.01	0.01
MnO	0.24	0.03	0.20	0.03	0.14	0.02	–	0.01
MgO	2.98	0.13	2.97	0.08	6.03	0.33	0.40	0.79
CaO	8.89	0.17	9.19	0.16	9.98	0.29	56.72	54.99
Na ₂ O	3.56	0.13	3.48	0.09	2.35	0.07	0.12	–
K ₂ O	2.05	0.10	2.05	0.03	4.01	0.19	–	–
P ₂ O ₅	0.29	0.04	0.34	0.04	0.70	0.09	0.03	0.02
Total	99.85		97.20		98.02		57.70	55.84
H ₂ O	–		2.23		2.00		0.15	–
CO ₂	–		–		–		44.93	44.11

^a Deegan et al. (2010)

^b Jolis et al. (2013)

^c Total Fe reported as FeO

^d 1 σ standard deviation; the number in parentheses represents the number of analyses for each sample

melt to be preserved. For a full account of starting materials, experiment preparation, experimental procedure, and rationale, see Deegan et al. (2010, 2011) and Jolis et al. (2013).

After the experiments were terminated, the capsules were mounted in epoxy resin, abraded to remove the capsule, and then polished to expose the surface of the experimental products at approximately mid-capsule depth. Back-scattered electron images of the polished surface were acquired using a JEOL JSM-6500F Field Emission Gun Scanning Electron Microscope (FEG-SEM) at INGV, Rome. As vesicle size varied over three orders of magnitude, image magnification was varied accordingly.

Image analysis and data treatment

The aim of this paper was to generate quantitative textural data to complement previous studies that employ quantification of the vesicle characteristics of volcanic rocks (e.g. Gardner et al. 1999; Shea et al. 2010; Heap et al. 2014). Low-magnification FEG-SEM images (as low as 15 \times magnification), covering the whole area of the exposed experimental charges, were used to measure the size of the largest vesicles. Conversely, images of increasingly higher magnification (up to 1500 \times) were used to measure smaller vesicles over progressively smaller areas. The number of images and the magnifications employed were modified for each sample to best capture its vesicle

size distribution. First, a visual inspection of the entire experiment at variable magnifications ensured the absence of clusters of small bubbles. This step was followed by taking the highest magnification images from randomly selected portions of the experiments. We then measured the area of each high-magnification image, including the areas of all chemical domains that contained vesicles, but excluded any area with remaining carbonate or calcite crystals. The area of all individual vesicles present in each high-magnification image was also measured separately using the ImageJ software (Abramoff et al. 2004). Vesicles attached to the edge of the capsule were included in the count, whereas those on the edge of an image were excluded. The Heywood diameter (equivalent diameter of a non-circular object) or Feret diameter (largest diameter obtainable from the area of a near-circular object) was calculated for each vesicle. Since these two measurements are equivalent for near-spherical vesicles, we will refer to the Heywood diameter only from here. These diameter values were ordered with the largest first and were assigned a cumulative number (C_n). To compensate for the smaller vesicle areas measured at high magnification, we corrected the number of bubbles to represent the whole capsule using a correction factor (F_{corr}). In this process, bubbles with a large area in the lower-magnification groups were removed from all successively higher-magnification groups to prevent smaller bubbles being extrapolated over areas that were occupied by larger bubbles.

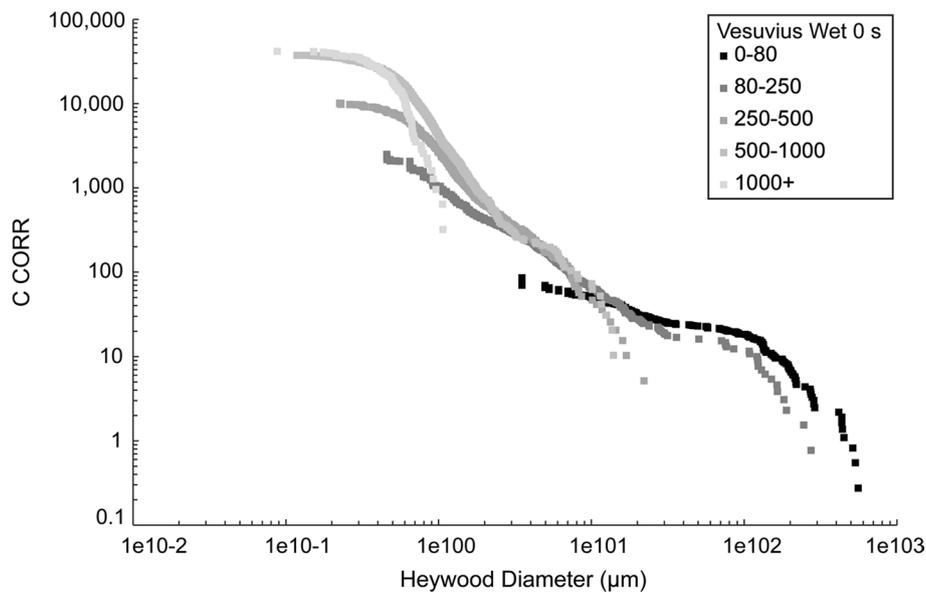


Fig. 1 Heywood diameter versus corrected cumulative number (C_{corr}) of vesicles for Vesuvius wet 0 s experiment. The cumulative vesicle correction extrapolates the number of vesicles measured over the area of the whole capsule. The data are firstly divided into respective magnification groups at which the vesicles were measured. Subsequently, the most representative magnification group for vesicles of

different diameters was identified. The 0–80 magnification group is most representative of, e.g. the largest vesicles, characterising those to ca. 16 μm in size, while the other curves are more representative of smaller sizes. A complete experimental curve can thus be reconstructed from the individual magnification curves (see Fig. 6c)

This correction factor (F_{corr}) and the cumulative correction (C_{corr}) were then calculated for each magnification class as follows:

$$F_{\text{corr}} = \text{Area}_{\text{capsule}} / \Sigma \text{Area}_{\text{image}} \quad (1)$$

$$C_{\text{corr}} = C_n \cdot F_{\text{corr}} \quad (2)$$

The cumulative correction value (C_{corr}) represents the theoretical number of bubbles present in the whole sample that are larger than any specific diameter (Fig. 1) and accounts for any over- or underestimation of bubble numbers produced in the measurement procedure. We used a Heywood diameter versus cumulative correction plot for each magnification group to identify where bubble measurements from each subsequently higher-magnification set appear. Since bubbles of almost all sizes were measured at least at two or three different magnifications, overlap between the magnification groups ensured the effectiveness of the area correction (Fig. 1). Finally, one single curve was created for each experimental charge from all magnification groups by using (1) representative magnifications to define the different sections of the curve and (2) favouring higher-magnification curves for smaller vesicle domains. Unfortunately, during polishing of the experimental charge 445-V7 (Vesuvius, 300 s), part of the sample was lost and the data are thus incomplete for this run. To compensate for this shortcoming, the missing area was not included in

higher-magnification area measurements, but was included in the whole sample area measurement. For extrapolation of vesicles over the whole sample area, we assumed that the lost area simply had the same vesicle arrangement as that observed in the remaining sample.

This method of vesicle size and number estimation employs a 2D surface and assumes that the largest diameter of each vesicle was intersected to derive an “ideal” diameter that forms the basis of our vesicle volume calculations. Our approach would introduce error into a 3D volume calculation as the largest diameter of each vesicle is not always intersected. Furthermore, as larger vesicles dominate the samples, the potential for smaller vesicles to be underrepresented in 3D is increased (e.g. Sahagian and Proussevitch 1998). The method employed here, however, produces an internally consistent data set for the experiments that is applicable for our purpose, but does not lend itself to accurate volume determinations.

Results

Textural observations

Here we provide a summary of the textural features recorded in the respective Merapi and Vesuvius experiments. This direct comparison of the two experimental

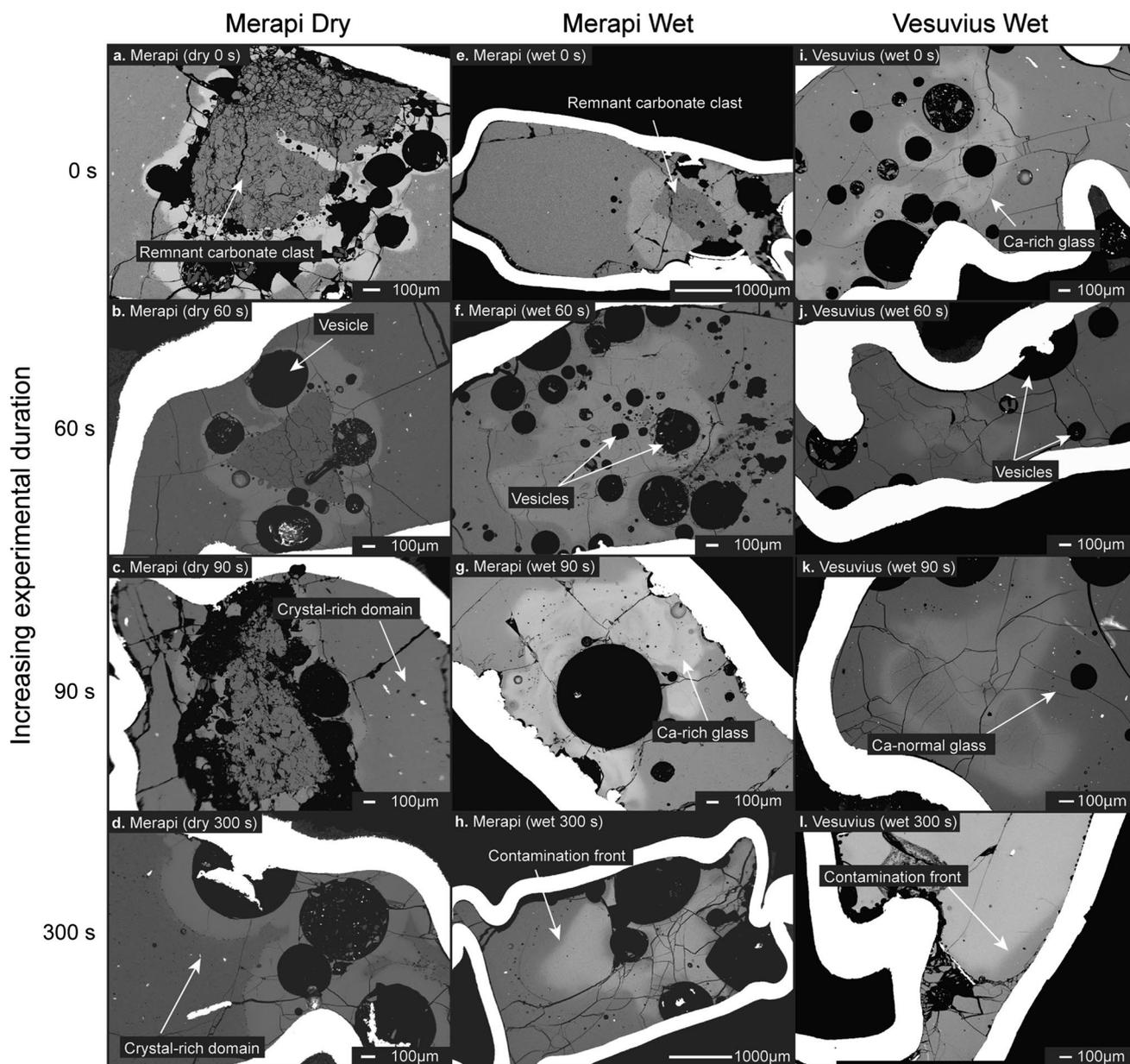


Fig. 2 Back-scatter electron images of textures and vesicle distributions produced in the three different experimental series and at four experimental durations. **a–d** Merapi dry series; **e–h** Merapi wet

series; and **i–l** Vesuvius wet series. Annotations show the remnant carbonate, the Ca-normal and crystal-rich domain, the Ca-rich glass, and, where appropriate, the contamination front

studies is intended to consolidate the available information in Deegan et al. (2010, 2011) and Jolis et al. (2013).

All three experimental series record chemical interaction between carbonate and melt as a result of partial or total assimilation of the carbonate (Fig. 2). Three compositional glass domains are identified: (1) a Ca-normal glass, compositionally similar to the starting material, (2) a Ca-rich glass, which arises due to breakdown of carbonate in the melt, and (3) a contamination front of varying extent that

usually marks the transition between the Ca-normal and Ca-rich domains (see Deegan et al. 2010 and Jolis et al. 2013 for full details). Assimilation of carbonate in these experiments is always accompanied by rapid vesiculation, and the size and spatial distribution of the vesicles are related to the stage of the reaction. Notably, in the dry series experiments, the Ca-normal domain contains crystal-rich regions, which affects melt viscosity (see discussion also).

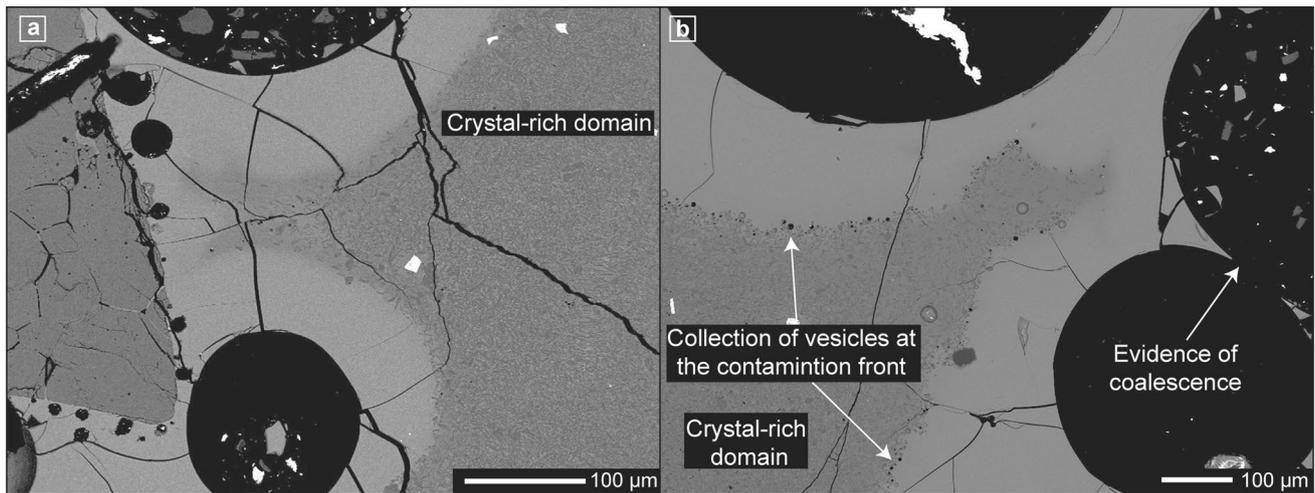


Fig. 3 Merapi dry 60 s experiment (a). Vesicles form at the interface between melt and carbonate (the reaction site). Smaller vesicles are generally located closer to the carbonate remnant. The contamination front, at the interface between the Ca-rich and the Ca-normal glasses, contours the outside edge of the larger vesicles. This suggests that bubble growth and bubble migration play a role in driving the contamination process in lower-viscosity systems. (b) Merapi dry

300 s experiment. Smaller vesicles collect at the contamination front between the crystal-rich domain and the Ca-normal glass. Two large vesicles can be seen to coalesce on the right-hand side of the image. Both quench crystals and smaller vesicles are present in the crystal-rich domain, while there are almost no small vesicles or crystals in the Ca-rich glass

Merapi dry series

In the Merapi dry series, the degree of carbonate assimilation increases with increasing run-time (i.e. from 0 to 90 s), with the most extensive carbonate incorporation into the melt at 300 s (Fig. 2a–d). At 0–90 s, the remaining carbonate is usually highly fractured and rounded and can show spallation of small fragments into the surrounding Ca-rich glass. The contamination front is well defined in all experiments of this series. Vesicles are primarily located in the Ca-rich glass with small vesicles frequently found near the carbonate, while larger vesicles occupy generally more distant regions in the capsules (e.g. Fig. 3a) and sometimes appear to concentrate at the contamination front. In fact, the contamination front frequently follows the profile of large bubbles (e.g. in the 60 and 300 s experiments; Fig. 2b, d), indicating that bubble growth and migration progressively drive the compositional interface away from the carbonate clast. Notably, a subset of smaller vesicles is also located in the crystal-rich domain (Fig. 2d).

Merapi wet series

In the Merapi wet series, the carbonate was no longer observed in the capsule by 90 s. The edges of the remaining carbonate in the 0 and 60 s runs are irregular, rounded, and also show spallation of carbonate clasts from the main carbonate fragment. Calcite crystals are present surrounding the remnant carbonate fragments in the 0 and 60 s

experiments and display a dendritic texture, which probably relates to rapid quenching of the experiments (Fig. 4a).

The extent of the Ca-rich glass generally increases at the expense of the Ca-normal one with increasing run-time (Fig. 2e–h). Ca-normal glass is observed enclosing areas of Ca-rich glass in the advanced time series, implying mingling of the two melts. The contamination front becomes wider and less defined at increasing experimental duration (cf. Deegan et al. 2010). Larger vesicles are located in both the Ca-rich and Ca-normal glasses, with an overall preference for the capsule margins. Smaller bubbles also occur in both glass types. However, in the Ca-normal glass, the smaller bubbles are arranged at a distance from the contamination front in semi-parallel bands, but maintain the shape of the contamination front (Fig. 4b). In addition, as in the Merapi dry experiments, the contamination front frequently follows the profile of large bubbles (e.g. 60 s experiment; Fig. 2f), highlighting the effect of vesicle growth and migration on the distribution of Ca-enriched (contaminated) melt domains.

Vesuvius wet series

In the Vesuvius wet series, the carbonate fragment was completely digested at 0 s. Although part of the 300 s experiment disintegrated during sample preparation, but because no carbonate was present at 0, 60, and 90 s, we assume that the carbonate was also fully dissolved in the 300 s run. In comparison with the Merapi series, the

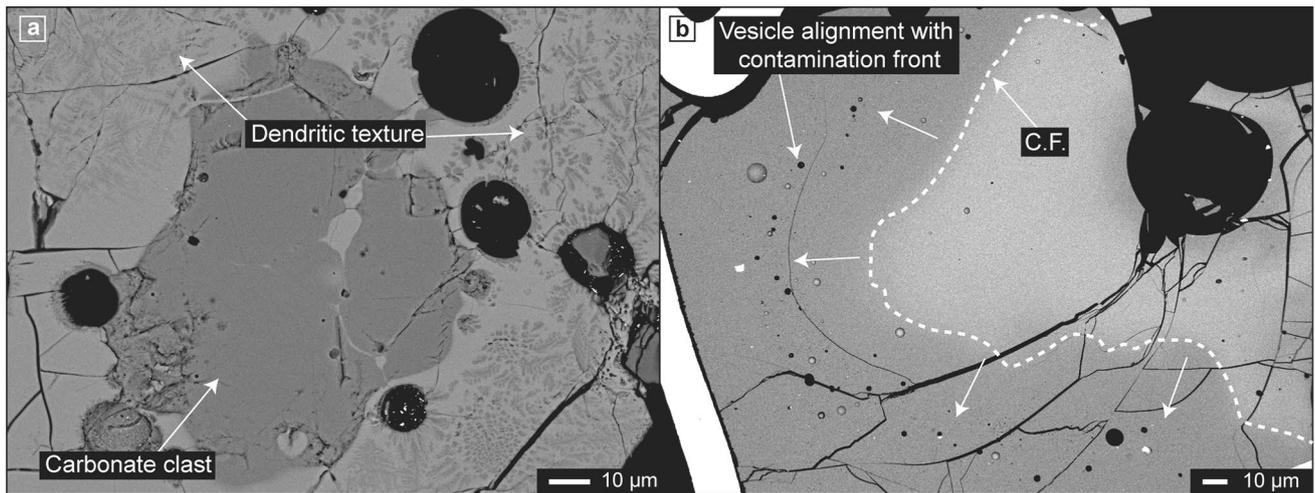


Fig. 4 Merapi wet 60 s experiment. **a** Example of dendritic texture present in some samples. This texture results from quenching of the Ca-rich melt that has digested calcite. Vesicles are located near the edge of the carbonate clast, some surrounded by microcrystals of calcite. **b** Merapi wet 300 s experiment. Vesicles of similar size are aligned parallel to the contamination front, but are located ahead of it in the Ca-normal glass. The lower viscosity of the Merapi wet experi-

ments relative to the Merapi dry ones allows volatiles to diffuse away from the reaction site more quickly. Here, all volatiles produced in one single vesiculation event seem to move away perpendicular to the reaction site at a similar rate. However, the viscosity is still relatively high in the Merapi wet series, and substantial mingling between the Ca-rich and Ca-normal melts is not yet observed

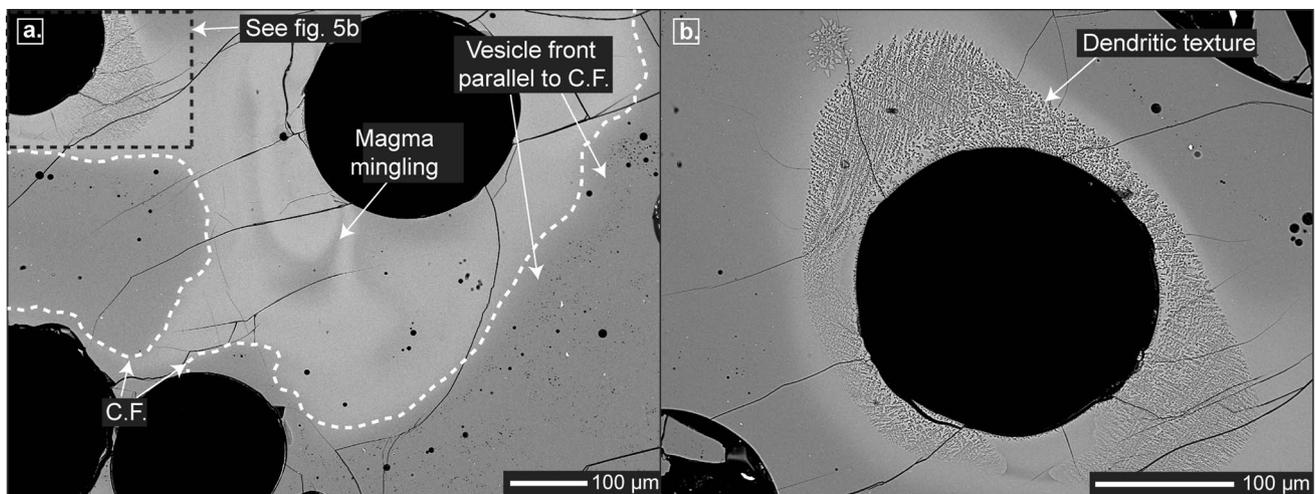


Fig. 5 Vesuvius wet 0 s experiment. **a** At 0 s, just as experimental temperature is reached, intermingling of Ca-rich and Ca-normal glass has already commenced, causing complex textures in comparison with the Merapi dry and wet series. Smaller vesicles in the Ca-normal glass align themselves along the contamination front (CF). Beyond the CF, i.e. further into the Ca-normal glass, smaller vesicles occur,

but do not form a regular pattern (e.g. Fig. 4b). The black stippled line around the dendritic textured area highlights the location of the higher-magnification image shown in (b). **b** The dendritic texture demonstrates a strong CaO enrichment, but is not associated with any solid carbonate remnants

Ca-rich and Ca-normal glasses are no longer as clearly defined and the contamination front is generally more diffuse, increasing in width as the experimental duration progresses (Jolis et al. 2013). Notably, the majority of vesicles at 0, 60, and 90 s are now in the Ca-normal glass. At 0 s, there are some larger bubbles located in the Ca-rich glass,

though the majority of smaller vesicles appear to be in the Ca-normal glass (Fig. 5a). At 300 s, Ca-rich glass is dominant in the experimental charge (see Jolis et al. 2013) and almost all vesicles are located in it. Calcite crystals are present in the 0, 60, and 90 s experiments and show a dendritic texture due to quenching (Fig. 5b).

Bubble nucleation and growth

Vesicle diameter and the cumulative vesicle number (C_{corr}) were obtained for each of the experimental series. In the Merapi dry series (Fig. 6 a), the 300 s experiment contains the largest amount of small bubbles ($<5 \mu\text{m}$), and overall, the size of the smallest bubbles decreases with increasing experimental duration (from $2.5 \mu\text{m}$ at 0 s to $0.2 \mu\text{m}$ at 300 s). Conversely, the size of the largest bubbles increases with increasing experiment duration (from $330 \mu\text{m}$ at 0 s to $880 \mu\text{m}$ at 300 s), and only the 90 s experiment does not follow these trends.

In the Merapi wet series (Fig. 6b), the 300 s experiment contains the largest amount of bubbles $<5 \mu\text{m}$ in size. The diameter of the smallest and largest bubbles increases from $0.1 \mu\text{m}$ at 60 s run to $0.4 \mu\text{m}$ at 300 s and from $\sim 135 \mu\text{m}$ at 0 s to $>1000 \mu\text{m}$ at 300 s, respectively. Both Merapi series show the total number of bubbles to increase with time, i.e. the 300 s experiments contain significantly more bubbles than the 0 s ones.

The Vesuvius wet series shows an analogous increase in maximum bubble size with time, but also a decreasing number of bubbles with increasing run-time (Fig. 6c). The smallest bubble size in the Vesuvius wet series ranges from $0.1 \mu\text{m}$ at 0 s to $0.3 \mu\text{m}$ at 300 s, whereas the largest bubble size at 60 s is ~ 470 and $\sim 610 \mu\text{m}$ at 300 s. The variation of the largest bubble size for the Vesuvius series (Fig. 6c) is therefore much less than in the two Merapi series (Fig. 6a, b). In this respect, vesicle diameter versus cumulative vesicle volume shows a general trend towards larger cumulative volumes associated with longer experimental duration for both of the Merapi series, as expected from the increasing degree of carbonate assimilation (Fig. 7a, b). In contrast, complete assimilation of carbonate is confirmed in the Vesuvius wet experiments by approximately equal cumulative bubble volumes in all experiments (Fig. 7c) and is consistent with the textural similarities to those Merapi experiments where extensive carbonate assimilation occurred (e.g. the 300 s Merapi experiments).

Discussion

Water and carbon dioxide solubility in Merapi and Vesuvius magmas

Given the presence of water and the addition of CO_2 to the experiments, it is important to consider volatile solubility in the starting materials to understand volatile exsolution and migration in the corresponding experiments (e.g. Dixon 1997; Vetere et al. 2011). For the Vesuvius wet series, the shoshonite melt can dissolve up to 7.9 wt% H_2O at $1250 \text{ }^\circ\text{C}$ and 0.4 GPa (Vetere et al. 2011). In the case of

Merapi, a basalt can contain 8.5 wt% H_2O at $1200 \text{ }^\circ\text{C}$ and 0.5 GPa (Newman and Lowenstern 2002) and an andesite up to 10 wt% H_2O at $1200 \text{ }^\circ\text{C}$ and 0.5 GPa (Botcharnikov et al. 2006). With respect to these data, all experimental glasses are H_2O under-saturated ($\text{H}_2\text{O} \leq 2 \text{ wt}\%$). However, the presence of CO_2 , whose quantity increases with progressive carbonate assimilation, will reduce water solubility in the melts (e.g. Holloway and Blank 1994; Botcharnikov et al. 2005; Deegan et al. 2010; Jolis et al. 2013). Using the volatile solubility model of Papale et al. (2006), we calculated H_2O – CO_2 solubilities for the Merapi wet experimental series. For the given temperature and pressure conditions ($1200 \text{ }^\circ\text{C}$ and 0.5 GPa), the Merapi wet experiments show a water solubility limit at 2.12 wt% and a CO_2 solubility of 0.36 wt%. As determined by Jolis et al. (2013), the Vesuvius wet experiments can dissolve up to 1.9 wt% H_2O and 0.34 wt% CO_2 . The water solubility for the Merapi wet and Vesuvius experiments in the presence of CO_2 is thus very close to the initial $\sim 2 \text{ wt}\%$ H_2O of the hydrous starting material, implying that when volatile saturation is reached, the evolving vapour phase is dominated by CO_2 . This can also be shown by employing the method of Deegan et al. (2010) and Jolis et al. (2013) to calculate the maximum amount of CO_2 produced from carbonate breakdown in the experiments. Taking the average masses of the magmatic starting material and the carbonate used in each experimental series, we calculate that 4.29 mg CO_2 (7.02 wt%), 4.28 mg CO_2 (7.0 wt%), and 2.9 mg CO_2 (4.75 wt%) can be produced for the Merapi dry, Merapi wet, and Vesuvius wet experiments, respectively. Carbon dioxide is thus more abundant in the experiments than H_2O ($\leq 2 \text{ wt}\%$) and represents the critical vapour phase in all series presented here.

Carbonate dissolution and the influence of melt viscosity

The rate of CaCO_3 breakdown in magma is controlled by the removal of CO_2 from the reaction site. Studies of natural systems such as the Hortavaer igneous complex (Norway) and Etna (Italy) have suggested that prolonged magma–carbonate interaction is permitted only if CO_2 is efficiently and continually removed from the reaction site (Barnes et al. 2003, 2005, 2009; Mollo et al. 2012; Heap et al. 2013). This is because decarbonation in a closed system would eventually slow down due to an excess of the reaction product. Although the removal of CO_2 in the experiments would eventually be constrained by the capsule, those experiments where the Ca-rich glass reaches the edge of the capsule show that bubbles were still growing larger with increased experiment duration, which suggests that outgassing is not yet limited by the capsules (cf. Fig. 2).

Fig. 6 Cumulative vesicle correction (C_{corr}) versus vesicle diameter (μm) for **a** Merapi dry, **b** Merapi wet, and **c** Vesuvius wet experiments. All four experiment durations (0, 60, 90, and 300 s) are plotted. The number of vesiculation events that can be identified increases from two in the Merapi wet 0 s to three in the Merapi wet 300 s. A similar sequence of events is also observed in the Vesuvius wet experiments. Vesicle coalescence can be identified where the number of bubbles in the smaller size fraction decreases with longer experimental duration (compare Fig. 7c)

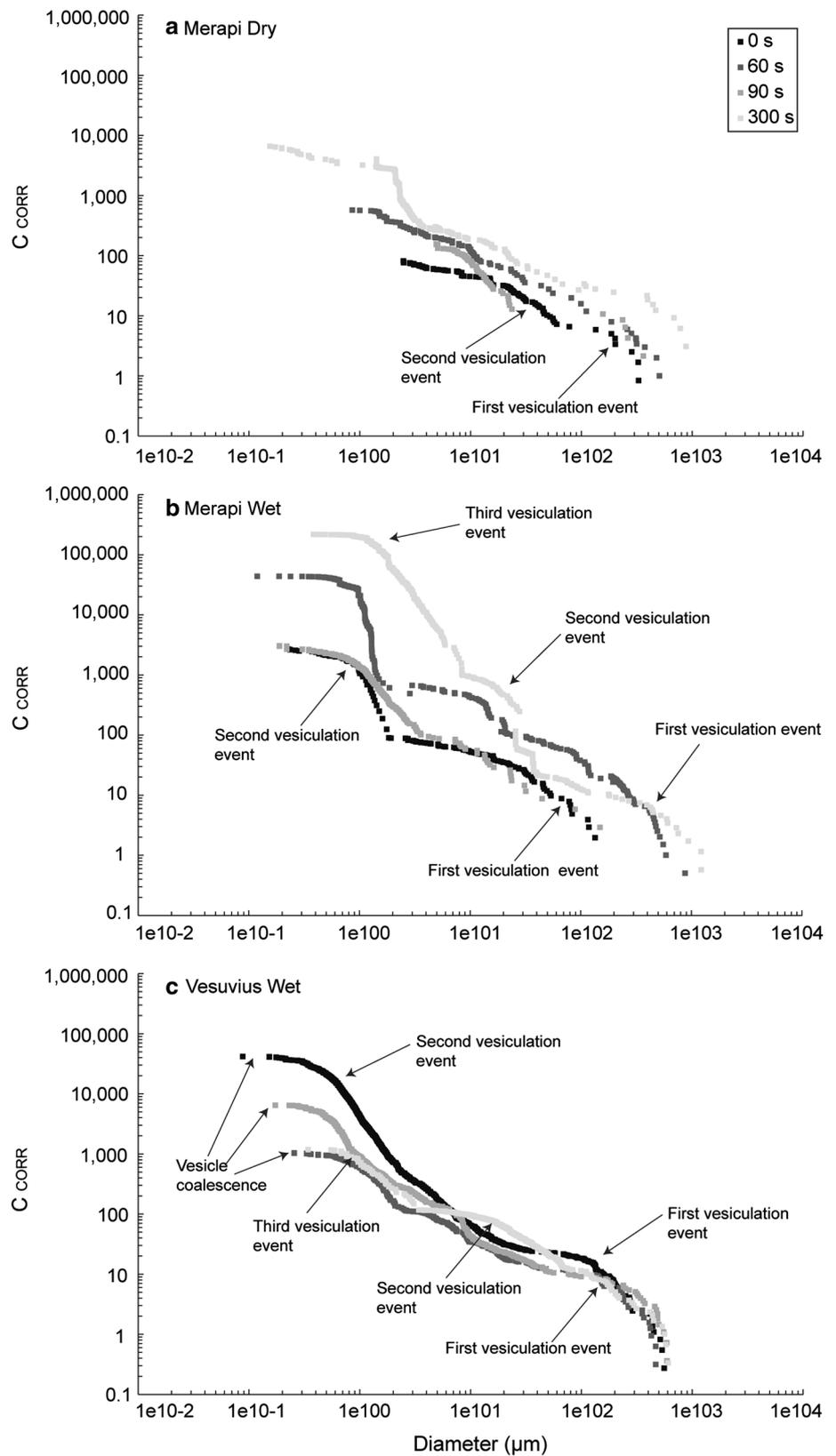
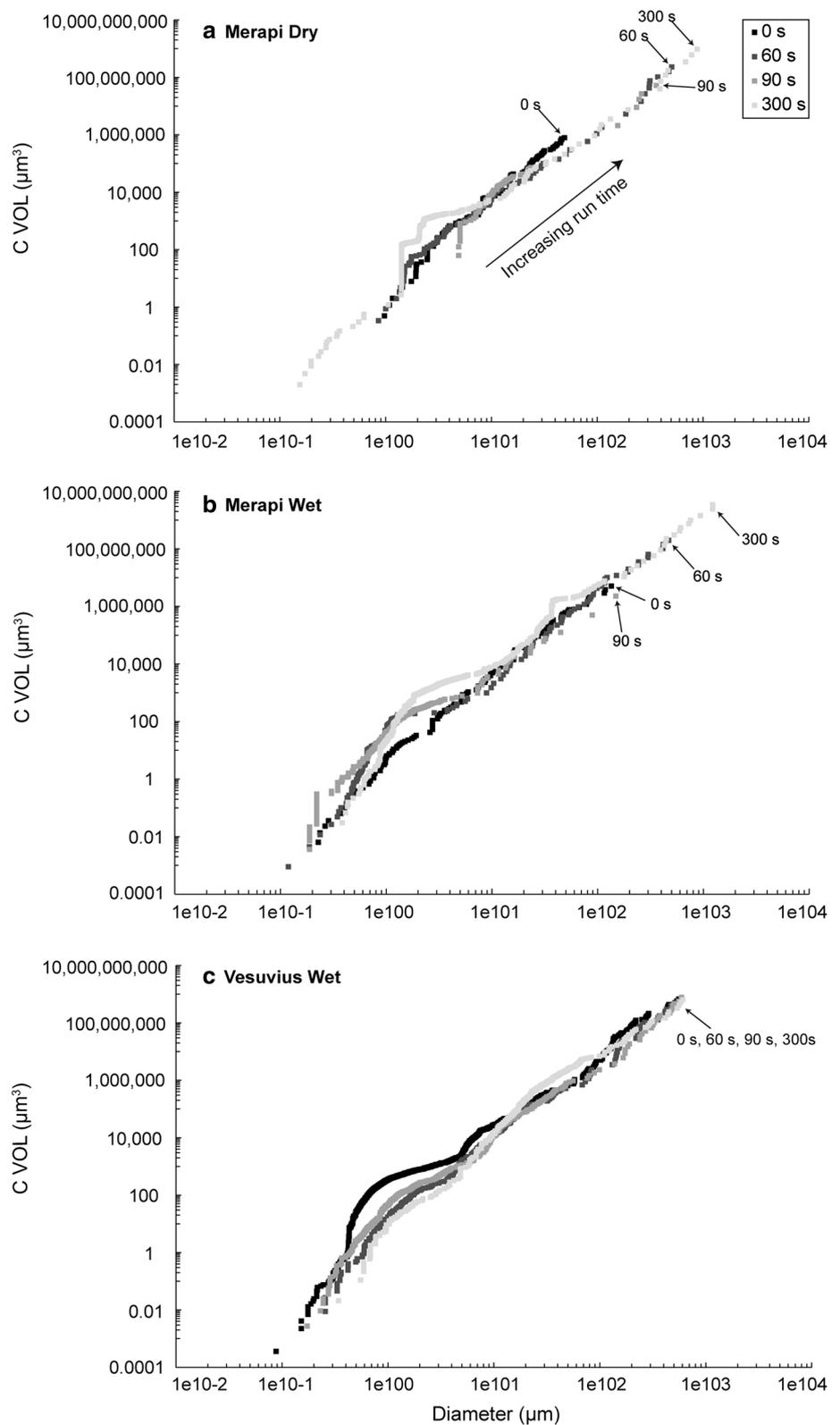


Fig. 7 Cumulative vesicle volume ($C_{vol} \mu\text{m}^3$) versus vesicle diameter (μm) for **a** Merapi dry, **b** Merapi wet, and **c** Vesuvius wet experiments. All four experiment durations (0, 60, 90, and 300 s) are plotted. For **a, b**, as experiment duration increases from 0 to 300 s, a general increase in the cumulative volume of bubbles is observed. This demonstrates that decarbonation has not reached completion in the shorter duration runs and that more CO_2 is exsolved in the longer duration experiments as the reaction proceeds. In **c**, the cumulative volume of vesicles is the same for all experimental durations. This indicates that by 0 s (i.e. just as experimental temperatures are reached), the decarbonation reaction is almost complete. As experimental duration increases for this series, homogenisation of the reaction products is the only change observed, i.e. no new bubbles are generated



In contrast, the experiments show that under the investigated conditions, carbonate dissolution began before the experimental target temperature was reached and thereafter

progressed with increasing experimental run-time. We note that the breakdown of carbonate is accelerated in the experiments over that of pure decarbonation because of reaction

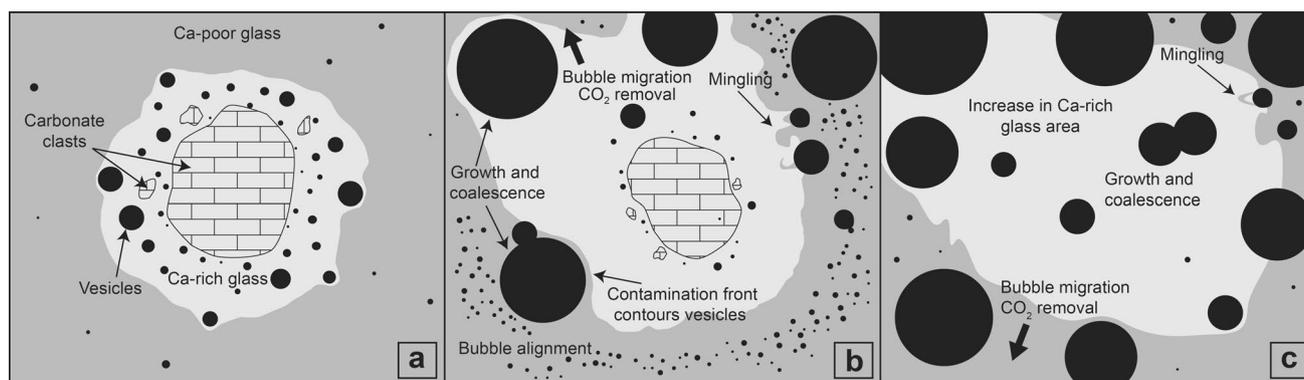


Fig. 8 Schematic diagram of carbonate dissolution processes. **a** Start of carbonate dissolution. Bubbles (*black filled circles*) nucleate mostly in the Ca-rich melt close to the dissolving carbonate fragment, although some appear also in the Ca-normal melt. **b** The carbonate clast decreases in size, and the Ca-rich melt increases in volume. Larger vesicles grow by diffusion and coalescence, pushing earlier bubbles and the contamination front away from the carbonate clast. Small vesicles parallel the contamination front, but are also found in Ca-normal melt where they formed as a result of volatile diffusion beyond the contamination front (Merapi wet and Vesuvius cases).

Stages **a, b** will be repeated multiple times as the carbonate clast dissolves in vesiculation pulses. The frequency and style of these pulses are controlled by the rate of removal of dissolved and exsolved CO_2 from the reaction site, which itself is controlled by melt viscosity. Once CO_2 has migrated from the reaction site, dissolution can proceed further, thus allowing for the next vesiculation pulse. **c** Eventually, the whole carbonate clast dissolves, and the majority of smaller bubbles coalesce into a small number of larger ones. Most vesicles have moved away from the reaction site, and bubble growth and migration promote significant mingling and mixing of melts

with magma. Bubble nucleation occurred heterogeneously, i.e. at the carbonate–melt interface, and homogeneously, i.e. within the melt. Heterogeneous nucleation also occurred on other available nucleation surfaces, such as microlites even at very low degrees of oversaturation (cf. Masotta et al. 2014) and probably explains the formation of small vesicles in crystal-rich domains (Fig. 2d). Homogeneous nucleation in certain melt domains, in turn, implies that volatiles were added (likely by diffusion) and that volatile saturation was then locally achieved (cf. Hurwitz and Navon 1994).

Indeed, Ca-rich glass and vesicles are present in all 0 s experiments, with the Vesuvius experiments representing the most extreme case wherein the carbonate was completely dissolved at 0 s. Comparing (1) the Merapi dry series, wherein carbonate assimilation required 300 s to reach an advanced stage and (2) the Merapi wet series, wherein the carbonate was largely dissolved by 90 s, it becomes apparent that compositionally controlled melt viscosity exerts a key control on the reaction rate.

For example, at 1200 °C, the Merapi dry, Merapi wet, and Vesuvius wet experiment starting materials (equivalent to the Ca-normal domains) have viscosities of $10^{2.7}$, $10^{1.4}$, and $10^{1.1}$ Pa.s, respectively (calculated after Giordano et al. 2008). If we take microcrystals in the Merapi dry series into account, we obtain a viscosity value of $10^{3.5}$ Pa.s. On the other hand, the Ca-rich melts have calculated viscosities between $10^{0.3}$ and $10^{0.2}$ Pa.s for the Merapi and Vesuvius series, respectively (assuming 25 wt% CaO), which is an order of magnitude lower than the results for the Ca-normal melts and confirms that the rate of carbonate dissolution is inversely related to melt viscosity.

In addition, the spatial distribution of vesicles in our time lapse experiments provides a record of bubble growth and migration. For all experiments, bubbles tend to move away from the carbonate clast as run-time proceeds, with greater vesicle mobility in lower-viscosity systems (Figs. 8, 9). Bubbles migrate away from the carbonate clast because a pressure and composition gradient is rapidly established in the experiments once the decarbonation reaction begins. The experiments quickly become oversaturated with CO_2 close to the reaction site, and bubble nucleation and migration become the main mass transport mechanisms. In the experiments, fluid pressure is greater at the reaction site than in the distal parts of the capsule, which causes bubbles to move away from the reaction site in a broadly radial fashion (Figs. 8, 9). Importantly, throughout the whole Vesuvius series, larger bubbles tend to be positioned preferentially at the capsule margins (Fig. 2i–l), reflecting fast volatile migration into the uncontaminated areas. Vesicle migration is sluggish in more viscous systems, however. For instance, in the Merapi dry series, where bubbles and Ca-rich melt remain relatively close to the carbonate clast (i.e. the reaction site) even at 300 s, virtually all bubbles are contained in the Ca-rich melt (Fig. 2d). Conversely, bubbles occur in both the Ca-rich and Ca-normal melt in the Merapi wet series, and there, the Ca-rich glass area is much wider. Notably, in the Merapi wet series, the boundary between Ca-rich and Ca-normal glasses follows the profile of large bubbles (Fig. 2f) or is aligned with many smaller bubbles (Figs. 2h, 4b).

In the Vesuvius wet 0 s experiment, mixing/physical mingling between the Ca-rich and the Ca-normal melts

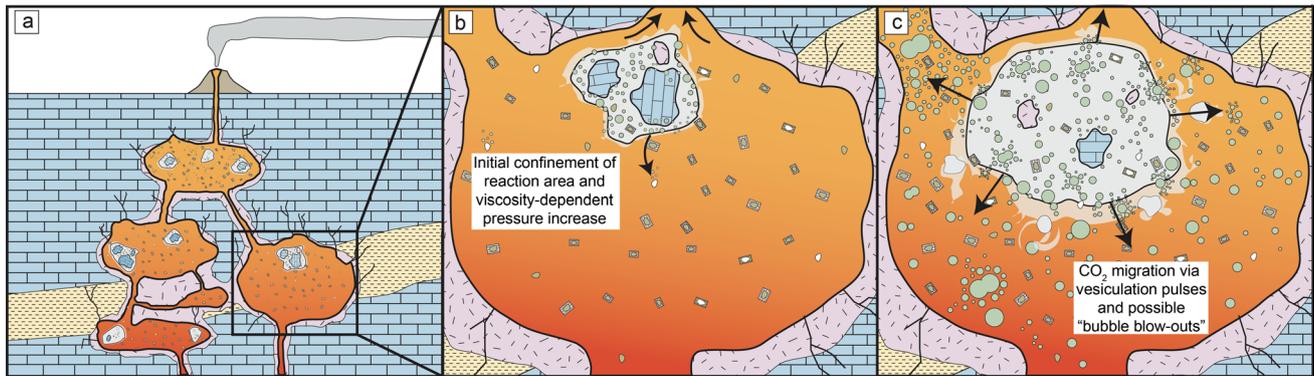


Fig. 9 Model depicting magma–carbonate interaction processes prior to an eruption. **a** The plumbing systems beneath both Vesuvius and Merapi consist of magma pockets distributed through the crust, which provide ample surface area for extensive magma–crust interaction. **b** Carbonate xenoliths are entrained into a slowly crystallising pocket of magma through, e.g. stoping, dyking, or earthquake activity. Reaction between carbonate blocks and the surrounding host melt produces CO₂ bubbles and a contaminated Ca-rich melt. The reaction area is initially localised and confined to the proximity of the carbonate. **c** As

the reaction continues, multiple vesiculation pulses cause CO₂ to penetrate a larger volume of host melt. Vesicle growth, coalescence, and migration thus promote mingling of Ca-rich and Ca-normal melts. If melt viscosity is relatively high (e.g. Merapi), volatiles are unable to freely escape and localised over-pressurisation can be achieved, resulting in an explosive “bubble blowout” event. Alternatively, if melt viscosity is relatively low (e.g. Vesuvius), efficient removal of CO₂ bubbles allows for a higher reaction rate, which results in more sustained volatile escape

is indicated by chemical gradients across the interface between the two glass domains. The interfaces between Ca-rich and Ca-normal melts roughly follow the shape of the bubbles, indicating that bubble migration facilitates melt mingling (Figs. 2, 5). We therefore suggest that the viscosity contrast between the Ca-rich and Ca-normal melts (see calculated values above) probably represents an initial hindrance for bubble migration and magma mixing, temporarily locking the bubbles in the Ca-rich melt as evident from the concentration of small bubbles along the contamination front in the Merapi dry series (Fig. 3b). In contrast, in the Merapi wet series, vesicles also grew in the Ca-normal zones. In this case, CO₂ pulsing from the carbonate resulted in CO₂ and bubble migration from the Ca-rich melt, as suggested by bubble trails in the Ca-normal glass that parallel the contamination front (e.g. Merapi wet 300 s; Fig. 4b). For the Vesuvius wet series, bubble nucleation also occurred in both Ca-rich and Ca-normal melts. The relatively low-viscosity contrast between these two melts appears to have promoted rapid bubble migration and efficient CO₂ diffusion into the Ca-normal domains (Fig. 5b).

Prolonged nucleation dominates the Merapi dry and wet runs where both small and large bubbles increase in number with time (Fig. 6a, b). The similar cumulative vesicle number for Merapi wet 60 s and Vesuvius wet 0 s experiments (Fig. 6), together with the increasing cumulative vesicle volume at longer run-times for both Merapi series (not observed in Vesuvius experiments; Fig. 7), demonstrates the different rates of assimilation-induced CO₂ exsolution. The slower assimilation rate causes prolonged bubble nucleation in both the Merapi dry and wet series, relative

to the Vesuvius wet series where nucleation is already exhausted as the carbonate is fully dissolved.

Fast digestion of carbonate in both Merapi and Vesuvius wet experiments hence causes sudden release of volatiles from the reaction site, but ambient melt viscosity controls their migration style. Bubble release likely occurs in a pulsatory fashion as multiple vesiculation events (or “bubble blow-outs”) are recognised from the sharp increases in vesicle number over a given diameter range in all of our experimental series (Fig. 6). In fact, the number of vesiculation pulses increases with experimental duration, and for example, the Vesuvius and Merapi wet experiments record two vesiculation pulses at 0 s, but three are seen at 300 s (see Fig. 6b, c). Although evidence of coalescence is found in all experimental series, in the form of, e.g. partially retracted bubble walls and films (e.g. Merapi dry series; Fig. 3b), coalescence becomes the dominant growth mechanism in the low-viscosity Vesuvius runs, where the overall number of bubbles (especially in the small size fraction) decreases with time (Fig. 6c).

Chemical gradients and “pile-up” of Ca and Sr between the Ca-rich and Ca-normal glass domains are found in both Merapi and Vesuvius hydrated experiments (Deegan et al. 2010; Jolis et al. 2013) and are of similar magnitude in all experiments (except for Vesuvius wet 300 s where homogenisation becomes increasingly relevant). The alignment of bubbles, both at and parallel to the contamination front, indicates a similar pile-up for CO₂ and underlines that volatile migration is influenced by viscosity through slowing down diffusive volatile migration and the movement of bubbles that have already formed.

Bubble formation affects both Ca-rich and Ca-normal domains, indicating that some CO₂ can diffuse through the rheological barrier at the melt interface(s) (Fig. 2). We also note that in low-viscosity systems, bubbles are efficiently transported away from the carbonate, causing a high frequency of bubble blowouts that can create a steady-state dissolution system. This may be analogous to an effervescent tablet in water that dissolves in very rapid pulses, usually almost too fast for us to see. In high-viscosity systems, in contrast, bubbles are initially locked in the melt close to the carbonate where they continue to grow, thus generating fluid overpressure and potentially episodic bubble blowout events. Initiation of a bubble blowout in such systems could be caused by several factors including (1) the bubble number density gradient between the contaminated and normal melt, (2) a progressive change of the chemical gradient, i.e. an increased tendency towards chemical homogenisation, or (3) external factors such as earthquakes, dyking, or magma replenishments (e.g. Deegan et al. 2011; Troll et al. 2012).

Implications for natural volcanic systems

The experiments represent laboratory simulations of magma–carbonate interaction, a process that has been shown to occur at both Merapi and Vesuvius volcanoes (Fulignati et al. 2004b; Chadwick et al. 2007; Iacono Marziano et al. 2009; Dallai et al. 2011; Borisova et al. 2013). The experiments are particularly relevant because they mirror features found in natural rocks. Micro-analyses of calc-silicate xenoliths from Merapi, for example, have been shown to contain extremely CaO-enriched, sometimes glassy domains (up to 61 wt% CaO, Deegan et al. 2010). At Vesuvius, Mg-rich olivine demonstrates elevated $\delta^{18}\text{O}$ values relative to typical mantle-derived mafic magmas due to carbonate-derived CO₂ fluxing in the roots of the volcanic system (Dallai et al. 2011). While the isothermal conditions of the experiments may not be fully applicable to assimilation at magma chamber walls, i.e. with a strong thermal contrast (cf. Del Moro et al. 2001; Spera and Bohrsen 2001), the experimental conditions replicate magmatic digestion of a stopped carbonate block or xenoliths from a “warm” chamber roof or wall (i.e. we assume the mass of the assimilated clast to be much less than that of the magma that provides heat to drive the reaction; Fig. 9). Specifically, we found that melt viscosity exerts a major control on the carbonate dissolution reaction by regulating the efficiency of bubble migration from the reaction site, which in turn determines how fast the carbonate can be digested. Low-viscosity systems such as Vesuvius will likely permit frequent vesiculation pulses where bubbles can migrate freely from the Ca-rich melt. Efficient migration of bubbles

from the dissolving carbonate can thus trigger more frequent new vesiculation pulses and promote higher assimilation rates (Fig. 9). This style of vesiculation would lead to steady-state CO₂ production for as long as the reactants are available. A potential outcome of this process would then be an increase in the buoyancy of Vesuvius magmas that would trigger rapid ascent and lead to sustained explosive eruptions (cf. Dallai et al. 2011).

For a more viscous magma (e.g. Merapi), bubble accumulation at the reaction site will progressively retard the carbonate dissolution reaction and the magma will locally behave as a closed system. If bubbles are not transported away, the reaction will simultaneously increase local fluid pressure until it terminates due to a lack of melt at the reaction site. If the confining pressure is exceeded during this process, however, a bubble blowout may be triggered in a broadly radial fashion from the reaction site, i.e. the carbonate (Figs. 8, 9), thus relieving localised fluid overpressure and allowing the decarbonation reaction to proceed again. This type of vesiculation would lead to an explosive, gas-driven eruption once the overpressure threshold is exceeded and may be repeated with successive CO₂ build-up events. This may apply to Merapi, especially the pre-2010 eruptions, as the latter was a “100-year event” (Surono et al. 2012) caused by a large contribution of fresh, low-viscosity basaltic magma from depth, rather than a dome-forming eruption as typical of the twentieth-century activity at Merapi. Notably, the bubble blowout mechanism seems to be also recorded in plutonic systems that underwent magma–carbonate interaction, such as the Hortavaer intrusive complex, Norway, where fragmented structures have been interpreted in this context (Barnes et al. 2009).

Regardless of the vesiculation style presented by a magmatic system undergoing magma–carbonate interaction, a fundamental finding from the experiments is that CO₂ fluxing from carbonate is coupled to magma viscosity. Where vesiculation becomes progressively sluggish, the bubble blowout mechanism described here can reignite vesiculation once the local fluid pressure threshold has been exceeded. In contrast, efficient CO₂ fluxing is permitted in relatively low-viscosity melts, which will allow carbonate assimilation to proceed unhindered and will likely lead to steady-state degassing.

The wider implication arising from these observations is that carbonate assimilation in a variety of different types of magmatic systems with differing melt properties can nevertheless contribute substantial quantities of CO₂ to the atmosphere via volcanism. It is thus imperative that carbonate-rich upper crustal lithologies beneath arcs are considered in mass balance models concerning Earth’s deep carbon cycle (cf. Johnston et al. 2011; Troll et al. 2012; Dasgupta 2013; Lee et al. 2013).

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